## **Beyond of the Poincaré Chern-Simons hypergravity, in (2+1) dimensions**

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## Introducción

We study the minimal coupling of Chern-Simons gravity based on the Maxwell symmetry with massless spin-5/2 gauge fields, in three dimensions. It is shown the simplest hyper-Maxwell superalgebra and its corresponding Chern-Simons gravity, which contains a massless spin-2 fields interacting with a massless Majorana spin-5/2 field. For certain combinations of osp(1/4) and sp(4) algebras, we find two different alternatives for hypersymmetric extensions of the Maxwell algebra, using the Inönü-Wigner contraction procedure. The hyper-Maxwell Chern-Simons actions was constructed for each case, which show interactions of non-propagating spin-4 fields with one or two spin-5/2 gauge fields.

## Desarrollo

We present two hypersymmetric extensions of the Maxwell algebra, which are obtained through the IW contraction procedure. The hypersymmetric extensions of the Maxwell CS gravity theory, was constructed by introducing fermionic generators which transform in an spin-3/2 irreducible representation of the Lorentz group. The sp(4) algebra is spanned by generators, sp(4) =  $\text{Span}\{T_a, T_{abc}\}$  and  $\text{osp}(1/4) = \text{Span}\{T_a, T_{abc}, G_{\alpha a}\}$ , where  $sl(2/R) \subseteq \text{sp}(4) \subseteq \text{osp}(1/4)$ . The simplest hypersymmetric extension of the Maxwell CS gravity theory, is spanned by the set

The simplest hypersymmetric extension of the Maxwell CS gravity theory, is spanned by the set  $\{J_a, P_a, Z_a, Q_{\alpha a}\}$ , where in addition to the usual commutators of Maxwell's algebra, the (anti-) commutation relations are:

 $[J_{a}, Q_{\alpha b}] = \frac{1}{2} (\Gamma_{a})^{\beta}_{.\alpha} Q_{\beta b} + \epsilon_{abc} Q^{..c}_{\alpha} , \quad \{Q_{\alpha a}, Q_{\alpha \beta}\} = -\frac{4}{3} \eta_{ab} Z_{c} (C\Gamma^{c})_{\alpha \beta} + \frac{5}{3} \epsilon_{abc} C_{\alpha \beta} Z^{c} + \frac{2}{3} Z_{(a|} (C\Gamma_{|b|})_{\alpha \beta}$ (1)

The  $T_a$  span the sl (2, R) subalgebra and stand the spin-2 generators,  $T_{abc}$  and  $G_{\alpha a}$  yields the spin-4 and spin-5/2 fields, respectively. The  $G_{\alpha a}$  are  $\Gamma$ -traceless vector-spinor ( $\Gamma^a G_a$ )<sub> $\alpha$ </sub> = 0 and  $T_{abc}$  are traceless ( $\eta_{ab}$ T<sub>abc</sub> = 0) and totally symmetric generators. The hyper-Maxwell algebras appear by considering the IW contractions (flat limit  $l \to \infty$ ) of the: Case I:  $osp(1/4) \otimes osp(1/4) \otimes sp(4)$ . This algebra, corresponds to a hyper-Maxwell algebra with spin-4 generators { $J_{abc}$ ,  $P_{abc}$ ,  $Z_{abc}$ }, and its hypersymmetric extension is characterized by two fermionic generators  $Q_{\alpha a}$  and  $\Sigma_{\alpha a}$  (required to satisfy the Jacobi Identities). The novel hypergravity can be seen as a Maxwellian generalization of the hyper-Poincaré gravity in presence of spin-4 gauge fields. Case II:  $osp(1/4) \otimes sp(4) \otimes sp(4)$ . This algebra, corresponds to a non-standard hyper-Maxwell (translational generators  $P_a$  are not expressed as bilinear expressions of fermionic generators  $Q_{\alpha a}$ ), where  $P_{abc}$  and  $J_{abc}$  transform in a spin-3 irreducible representation of the Lorentz group. Using a Lie algebra expansion method, a generalization of hyper-Poincaré and hyper-AdS gravities can be found.

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## Referencias

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